3- D coherent dispersion effect with space charge

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- 1. Combined effect of space charge and dispersion
- 2. Space-charge-modified dispersion function
- 3. Coherent dispersion mode

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- 4. Beam instabiltiies with dispersion
- 5. Splitting effect in 3-D bunched beam

Combined effect of space charge and dispersion

Two fundamental effects of accelerators

 Beam dynamics in high-intensity circular accelerator are subject to combined effect of space charge and dispersion

Combined effect of space charge and dispersion

- \triangleright Dispersion effect (only) \Box increase horizontal beam size
	- \checkmark Increase beam size
- \triangleright Space charge (only)
- \checkmark modify the lattice functions
- \checkmark envelope modes/ instabilities

 \triangleright Dispersion effect + space charge

- \checkmark Increase horizontal beam size
- \checkmark modify the lattice functions
- coherent **dispersion** (envelope) mode
- \checkmark induce coherent beam instabilities

Tools: combined rms envelope approach

 \triangleright Starting potint Hamiltonian for single particle with s.c. and disp.

$$
H = \frac{1}{2} (p_x^2 + p_x^2) + \frac{\kappa_{x0}(s)}{2} x^2 + \frac{\kappa_{y0}(s)}{2} y^2 + \frac{m^2 c^4}{E_0^2} \delta^2 - \frac{x}{\rho(s)} \delta + V_{sc}(x, y, s)
$$

Using Venturnin-Reiser and Lee-Okmoto theory

$$
x = \overline{x} + \overline{S}D_x
$$

$$
x' = \overline{x}' + \overline{S}D_x
$$

$$
y = \overline{y}
$$

$$
y' = \overline{y}'
$$

Betatron coordinate with s.c.

Off-momtenum coordinate with s.c. M. Venturini and M. Reiser, Phys. Rev. Lett. 81, 96 (1998)

Tiglildi di liulgo.

S. Y. Lee and H. Okamoto, Phys. Rev. Lett. 80, 5133 (1998) Following V-R and L-O theory, beam coherent motion

betatron montion

dispersion motion

How to seperate/identify the two parts in the total beam size?

• By defining the space charge-modified dispersion

$$
\frac{d^2 D_x}{ds^2} + \left[K_{x0}(s) - \frac{K_{sc}}{2X(X+Y)}\right]D_x = \frac{1}{\rho(s)}
$$

Space charge depression

X,Y : total beam size

Space- charge modified dispersion function

$$
\frac{d^2 D_x}{ds^2} + \left[\kappa_{x0}(s) - \frac{K_{sc}}{2X(X+Y)}\right]D_x = \frac{1}{\rho(s)}
$$

 \triangleright Ansatz 1 $\langle \overline{x} \delta \rangle = 0$

• Second order moment $X^2 = (\overline{x} + \delta\!D_x)^2 = \overline{x}^2 + 2D_x \big\langle \overline{x} \delta \big\rangle + \big(\delta\!D_x \big)^2 = \overline{x}^2 + \big(\delta\!D_x \big)^2$ $=\sigma_x^2+\sigma_p^2D_x^2$ of the total beam size

Ansatz 2

• rms equivalence of Frank Sacherer holds with dispersion

$$
\left\langle x \frac{\partial V_{sc}}{\partial x} \right\rangle = -\frac{K_{sc}}{2} \frac{X}{X+Y}
$$

Space- charge modified dispersion function

 \triangleright With the definition of Dx, the parts are independent with each other.

$$
X^2 = (\overline{x} + \delta D_x)^2 = \overline{x}^2 + (\delta D_x)^2 = \sigma_x^2 + \sigma_p^2 D_x^2
$$

Envelope rms equation with dispersion

 \triangleright By using the coordinate transformation

$$
\varepsilon_{dx} = \sqrt{\langle \overline{x}^2 \rangle \langle \overline{x}'^2 \rangle - \langle \overline{x}' \overline{x} \rangle^2}
$$

 \triangleright The rms envleope equaion combined with dispersion

$$
\frac{d^2 \sigma_x}{ds^2} + \left[\kappa_{x0}(s) - \frac{K_{sc}}{2X(X+Y)}\right] \sigma_x - \frac{\varepsilon_{dx}^2}{\sigma_x} = 0
$$

$$
\frac{d^2 \sigma_y}{ds^2} + \left[\kappa_{y0}(s) - \frac{K_{sc}}{2Y(X+Y)}\right] \sigma_y - \frac{\varepsilon_{dy}^2}{\sigma_y} = 0
$$

$$
\frac{d^2 D_x}{ds^2} + \left[\kappa_{x0}(s) - \frac{K_{sc}}{2X(X+Y)}\right]D_x = \frac{1}{\rho(s)}
$$

Pioneer works on the genelized envelope equation set

J. A. Holmes, V. V. Danilov, J. D. Galambos, D. Jeon, and D. K. Olsen, Phys. Rev. ST Accel. Beams 2, 114202 (1999).

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CFD model

- **C**onstant **F**ocusing channel with **D**ispersion (smooth approximation model)
	- e.g. bending solenoid focusing structure, the matched beam sizes are given by

Characterisitcs of space- charge modified dispersion

Two characterisitcs of the s.c.-modified dispersio can be shown by CFD

(1) For matched beams

- As beam current increasing, the beam experiences two stages
	- 1. X/Y increasing: dispersion dominated stage

2. X/Y decreasing: space-charge dominated stage

Tune depression

$$
\eta_{x,y} = \frac{k_{0,x,y} - k_{x,y}}{k_{0,x,y}}
$$

Characterisitcs of space- charge modified dispersion

(2) For mismatch oscillations

 Coherent oscillaton of mismatched beams can be analyzed via perturvation on the combined envelope equation

 \triangleright Here coefficients a₁ to a₅ are functions of matched beams and emittance

Characterisitcs of Space- charge modified dispersion

(2) For mismatch oscillations

Envelope modes ϕ_1 , ϕ_2

and dispersion mode ϕ_3

can be identified from the coefficient matrix of a1 to a5

The limit of the three modes are equal to that in the case without dispersion

Pioneer works on dispersion mode

M. Ikegami, S. Machida, and T. Uesugi, Phys. Rev. ST Accel. Beams 2, 124201 (1999)

Coherent oscillations and dispersion mode

 \triangleright PIC simulations are in agreement with the numerical envelope approach

 \triangleright The frequency of the three modes can be obtained from the beam spectrum of the second order moments

$$
\langle x^2 \rangle
$$
 and $\langle y^2 \rangle$

Coherent beam instabilities with dispersion

 \triangleright For alternating focusing structure, the mismatched oscillaiton can drive instabilites

 \triangleright The most well-known is the envelope instability (second order even instability)

 \triangleright Dispersion+space charge \rightarrow modify the envelope instability drive the dispersion instability

 η_x

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Coherent beam instabilities with dispersion

- \triangleright Dispersion-induced instability
	- Dispersion mode is on resonacen with periodic focusing structure
	- criteria: phase advance > 120 degree
- \triangleright Compared with: Envelope instability
	- Envelope modes are resonant with periodic focusing structure
	- Criteria: phase advance > 90 degree

3- D bunch case

For 3D bunches, dispersion mode will be split due to the synchrotron motion

Analysis:

$$
\int \xi'' + a_0 \xi + a_1 \eta + a_2 d_x = 0
$$

\n
$$
\eta'' + a_1 \xi + a_3 \eta + a_4 d_x = 0
$$

\n
$$
d_x'' + \frac{a_2}{\sigma_p^2} \xi + \frac{a_4}{\sigma_p^2} \eta + a_5 d_x = 0
$$

$$
X^{2} = \langle x^{2} \rangle = (\sigma_{x} + \xi)^{2} + \left((D_{x} + d_{x})^{2} \sigma_{p}^{2} \right)
$$

$$
Y^2 = \langle y^2 \rangle = (\sigma_y + \eta)^2
$$

3- D bunch case

\triangleright For 2-D coasting beams

$$
\sigma_p = \text{const} \qquad (D_{x0} + d_x)\sigma_{p0}
$$

 \triangleright For 3-D bunched beams

$$
\sigma_p = \sigma_{p0} + \Delta \sigma_p(t) \qquad (D_{x0} + d_x)[\sigma_{p0} + \Delta \sigma_p(t)]
$$

 \triangleright The width of the split "gap" depends on the synchrotron frequency

3 - D bunch case

- \triangleright Sidebands appear around the envleope modes
- \triangleright In the presence of space charge, the split of dispersion mode is coupled to the envelope modes
- \triangleright space-charge coupling between

 $(D_{x0} + d_x)[\sigma_{p0} + \Delta \sigma_p(t)]$

and

ξ,^η

- Combined effect of space charge and dispersion has been investigated by using the envelope equation set including dispersion
- \triangleright For 2-D coasting beams,
	- Envelope instabilies are modified
	- Disperison-induced instability
- For 3-D bunced beams, dispersion mode is split because of synchrotron motion

In the future

 \triangleright Chromoticity effect on the dispersion mode

